

Diffusion in a 3D SUSY σ - model

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Outline of Talk

A) Motivation: Study time evolution of quantum particle in a random environment using **Statistical Mechanics**.

Advantage: Saddle point analysis + Symmetries

B) Spectral properties algebraically equivalent to correlations in a Supersymmetric Statistical Mechanics Model:

Hyperbolic, $SU(1, 1|2)$ symmetry,

Spin = 4x4 super-matrix,

Partition function $\equiv 1$.

C) Result: Diffusion for a **simpler** 3D SUSY model:

Model due to Zirnbauer, has hyperbolic supersymmetry.

Expected to reflect a Localization-Diffusion transition.

Equivalent to random walk in a *correlated* random environment.

D) Key Ideas of proof:

Estimate fluctuations of the environment
using SUSY Ward Identities.

Estimates on **nonuniformly** elliptic equations.

E) Relation to Edge Reinforced Random Walk?

This is *history* dependent random walk which favors edges it has visited in the past.

It is also equivalent to a random walk in a *correlated* random environment (Diaconis).

Localization in 1D + partial results in 2D (Merkl, Rolles) using a Mermin-Wagner type argument to estimate fluctuations of the environment.

Quantum Green's Function

Let \mathbf{H} denote a *random* band matrix or Schrödinger operator.

$$G_\epsilon(E; x, y) \equiv (\mathbf{H} - E + i\epsilon)^{-1}(x, y).$$

Q-Diffusion Conjecture in 3D:

$$\langle |G_\epsilon(E; x, y)|^2 \rangle \cong \frac{\rho(E)}{-D\Delta + \epsilon}(x, y) \approx C |x - y|^{-1}$$

$\rho(E) = \text{Im} \langle G_\epsilon(E; 0, 0) \rangle =$ density of states

$D = D(E) > 0$ is diffusion constant.

Localization :

$$\langle |G_\epsilon(E; x, y)|^2 \rangle \cong \epsilon^{-1} e^{-|x-y|/\ell}$$

SUSY Hyperbolic Sigma model

Sigma constraint: $-z_j^2 + x_j^2 + y_j^2 + 2\bar{\psi}_j\psi_j = -1$.

Action \mathbf{A} in *horospherical* coordinates:

$t_j, s_j \in \mathbb{R}$, and $\bar{\psi}_j, \psi_j$ Grassmann

$$\mathbf{A}_\epsilon(\mathbf{t}, \mathbf{s}, \psi) \equiv \sum_{j \sim j'} (\cosh(\mathbf{t}_j - \mathbf{t}_{j'}) - 1) + \frac{\epsilon}{\beta} \sum_j \cosh \mathbf{t}_j$$

$$+ \sum_{j \sim j'} e^{(\mathbf{t}_j + \mathbf{t}_{j'})} \left[\frac{1}{2} (\mathbf{s}_j - \mathbf{s}_{j'})^2 + (\bar{\psi}_j - \bar{\psi}_{j'}) (\psi_j - \psi_{j'}) \right]$$

$$Z = \int e^{-\beta \mathbf{A}_\epsilon(\mathbf{t}, \mathbf{s}, \psi)} \prod_j e^{-\mathbf{t}_j} dt_j ds_j d\psi_j d\bar{\psi}_j \equiv 1$$

Elliptic generator \mathbf{D} with random conductances

For $t_j \in \mathbb{R}$, define $\mathbf{D}(\mathbf{t})$ via the quadratic form:

$$(f, \mathbf{D}_\beta(t) f) = \beta \sum_{j \sim j'} e^{t_j + t_{j'}} (f(j) - f(j'))^2.$$

\mathbf{D} is the generator of *random walk in the random environment* of the t field. \mathbf{D} is **not** uniformly elliptic.

If $e^{(t_j + t_{j'})} \sim 1$ then there is **diffusion**.

However, if $\langle e^{t_j/2} \rangle_{\beta, \epsilon} \ll 1$, then conductance goes to zero and **localization** occurs. Known in 1D and expected in 2D

Relation to Quantum-Green's Function

$$\langle |G_{\epsilon}(E; x, y)|^2 \rangle_{RM} \cong \langle s_x e^{t_x} s_y e^{t_y} \rangle_{SUSY}(\beta, \epsilon)$$

$$= \langle [\beta \mathbf{D}(t) + \epsilon e^t]^{-1}(x, y) e^{t_x + t_y} \rangle_{SUSY}(\beta, \epsilon)$$

$$= \langle [-\beta \Delta + V(t) + \epsilon e^{-t}]^{-1}(x, y) \rangle_{SUSY}(\beta, \epsilon)$$

$\beta \sim$ density of states $\rho(E)$, times band width.

Main Theorem: "Diffusion" in 3D (Di-Sp-Zi):

For β **large**, local conductance $e^{t_j+t_{j'}} \sim 1$, hence :

$$\begin{aligned} & \langle e^{t_0+t_x} [\beta \mathbf{D}(t) + \epsilon e^t]^{-1}(0, x) \rangle_{SUSY}(\beta) \\ & \approx (-\beta \Delta + \epsilon)^{-1}(0, x) \approx C|x|^{-1} \end{aligned}$$

So, the evolution is "**Diffusive**".

Conjecture: For **small** β , **localization** occurs in 3D:

$\langle e^{t_j/2} \rangle \rightarrow 0$ as $\epsilon \rightarrow 0$, conductance vanishes.

Ward Identities are Key to Proof

To estimate the conductance, we first bound the *fluctuations* of the t field,

$$\langle \cosh^m(t_0 - t_\ell) \rangle (\beta, \epsilon) \leq \text{Const.}$$

by induction on $|\ell|$ and Ward identities:

Example: For $\ell \in \mathbb{Z}^d$, let

$$F_\ell = \cosh(t_0 - t_\ell) + e^{t_0 + t_\ell} \left[\frac{1}{2}(s_0 - s_\ell)^2 + (\bar{\psi}_0 - \bar{\psi}_\ell)(\psi_0 - \psi_\ell) \right]$$

Then $\langle F_\ell^m \rangle = 1$ all $m \geq 1$.

After integrating over the ψ and s variables

The Ward identity $\langle F_\ell^m \rangle = 1$ yields:

$$\langle \cosh^m(t_0 - t_\ell) (1 - \frac{m}{\beta} G_\ell(t)) \rangle = 1$$

where G_ℓ is the Green's function of $D(t)$.

If $0 \leq G_\ell(t) \leq \mathbf{C}$, then for $\mathbf{C}m < \beta$

$$\langle \cosh^m(t_0 - t_\ell) \rangle \leq (1 - \frac{\mathbf{C}m}{\beta})^{-1}$$

Thus t - fluctuations at scale ℓ are small for large β .

Conclusions and open Problems

Quantum Dynamics \approx SUSY hyperbolic model = RWRE.

In our Hyperbolic SUSY model, localization and diffusion are revealed using a simple **saddle** analysis. But **fluctuations** need to be **estimated**. We do this in 3D for large β .

Once the Ward identities are established, **analysis is classical**:
Induction on length scales + estimates on elliptic Greens functions.

Problem: Prove Localization in 2D and Phase Transition in 3D.

Happy Birthday Yoseph!