

Matrix Normal Forms

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I. INTRODUCTION TO THE PROBLEM

A. Example: Gauss-Jordan elimination

Problem: solve a system $A\vec{x} = \vec{b}$ of m equations in n unknowns, with coefficients in a field \mathbb{F} . Here

$$A = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ a_{21} & a_{22} & \cdots & a_{2n} \\ \vdots & \vdots & & \vdots \\ a_{m1} & a_{m2} & \cdots & a_{mn} \end{bmatrix} \quad \vec{x} = \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ \vdots \\ x_n \end{bmatrix} \quad \vec{b} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_m \end{bmatrix}$$

Method of solution: *row reduction* of the augmented matrix $[A \mid \vec{b}]$:

$$B = [A \mid \vec{b}] = \left[\begin{array}{cccc|c} a_{11} & a_{12} & \cdots & a_{1n} & b_1 \\ a_{21} & a_{22} & \cdots & a_{2n} & b_2 \\ \vdots & \vdots & & \vdots & \vdots \\ a_{m1} & a_{m2} & \cdots & a_{mn} & b_m \end{array} \right] \xrightarrow[\text{operations}]{\text{elementary row}} R,$$

to obtain a matrix R in *reduced row-echelon form (RREF)*.

The relation between B and R is called *row equivalence*. Alternatively, R and B are row equivalent if $R = PB$ with P $m \times m$ invertible.

A matrix in RREF looks like this (● represents any number):

$$\begin{bmatrix} 0 & \mathbf{1} & \bullet & \bullet & 0 & 0 & \bullet & 0 & \bullet & \cdots \\ 0 & 0 & 0 & 0 & \mathbf{1} & 0 & \bullet & 0 & \bullet & \cdots \\ 0 & 0 & 0 & 0 & 0 & \mathbf{1} & \bullet & 0 & \bullet & \cdots \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \mathbf{1} & \bullet & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \cdots \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots \end{bmatrix}$$

More precisely, a matrix R is in reduced row-echelon form if

- All nonzero rows lie above any zero rows;
- The first nonzero entry (the *pivot*) in any nonzero row is a 1 (this distinguishes the *pivot columns* $j_1 < \cdots < j_r$);
- Each pivot lies to the right of the pivot in the row above it;
- All matrix entries above a pivot are zero.

Remark: the number of nonzero rows in R is the *rank* of R —or of B , since row equivalent matrices have the same rank.

B. A normal form theorem

Theorem: Every $m \times n$ matrix is row equivalent to a unique matrix in reduced row-echelon form.

This is the standard setup in which *normal* or *canonical forms* for matrices arise. We have:

- An equivalence relation on a set \mathcal{S} of matrices;
- A special class \mathcal{N} of matrices within the set, the *matrices in normal form*, such that every matrix in \mathcal{S} is equivalent to a unique (or essentially unique) matrix in \mathcal{N} .

Thus the RREF is a normal form for rectangular matrices under row equivalence.

C. Finish discussion of linear equations

Consider again $A\vec{x} = \vec{b}$. Suppose under row reduction

$$[A \mid \vec{b}] \longrightarrow [S \mid \vec{c}] = \left[\begin{array}{cccccccc|cc} x_1 & x_2 & x_3 & x_4 & x_5 & x_6 & x_7 & & x_{n-1} & x_n & & \\ \hline 0 & \mathbf{1} & \bullet & \bullet & 0 & 0 & \bullet & \cdots & 0 & \bullet & & c_1 \\ 0 & 0 & 0 & 0 & \mathbf{1} & 0 & \bullet & \cdots & 0 & \bullet & & c_2 \\ 0 & 0 & 0 & 0 & 0 & \mathbf{1} & \bullet & \cdots & 0 & \bullet & & c_3 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & & \vdots \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots & \mathbf{1} & \bullet & & c_r \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 & & c_{r+1} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 & & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & & \vdots \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & \cdots & 0 & 0 & & 0 \end{array} \right]$$

The equations $S\vec{x} = \vec{c}$ are equivalent to the original $A\vec{x} = \vec{b}$.

Here c_{r+1} is 0 or 1. There is no solution if $c_{r+1} = 1$.

If $c_{r+1} = 0$, solve for the variables $x_2, x_5, x_6, \dots, x_{n-1}$ corresponding to the pivot columns; the other variables $x_1, x_3, x_4, x_7, \dots, x_n$ are free parameters.

D. Linear equations in abstract vector spaces

V, W finite dimensional vector spaces over \mathbb{F} , $L : V \rightarrow W$ linear.

Problem: Given $\mathbf{b} \in W$ find $\mathbf{x} \in V$ with $L\mathbf{x} = \mathbf{b}$.

To solve this problem, introduce:

- *Ordered bases:* $(\mathbf{v}_1, \dots, \mathbf{v}_n)$ for V and $(\mathbf{w}_1, \dots, \mathbf{w}_m)$ for W
- *Coordinates in each vector space:* If $\mathbf{x} \in V$ has the expansion $\mathbf{x} = \sum_{i=1}^n x_i \mathbf{v}_i$, and $\mathbf{b} \in W$ the expansion $\mathbf{b} = \sum_{i=1}^m b_i \mathbf{w}_i$, then \mathbf{x} and \mathbf{b} have *coordinate vectors*

$$\vec{x} = \begin{bmatrix} x_1 \\ \vdots \\ x_n \end{bmatrix} \quad \text{and} \quad \vec{b} = \begin{bmatrix} b_1 \\ \vdots \\ b_m \end{bmatrix}.$$

- *The matrix A of the linear operator L :*

$$L\mathbf{v}_i = \sum_{j=1}^m a_{ji} \mathbf{w}_j, \quad i = 1, \dots, n$$

- Then $L\mathbf{x} = \mathbf{b} \leftrightarrow A\vec{x} = \vec{b}$ and we solve the problem as above.

Notation: It is convenient to write the bases as column vectors whose components are abstract vectors: $\vec{v} = \begin{bmatrix} \mathbf{v}_1 \\ \vdots \\ \mathbf{v}_n \end{bmatrix}$, $\vec{w} = \begin{bmatrix} \mathbf{w}_1 \\ \vdots \\ \mathbf{w}_m \end{bmatrix}$.

Then $\mathbf{x} = \vec{v}^T \vec{x}$, $\mathbf{b} = \vec{w}^T \vec{b}$ and $L\mathbf{x} = \vec{w}^T A \vec{x}$.

Basis change. Consider different choices of bases, say $(\mathbf{v}'_1, \dots, \mathbf{v}'_n)$ and $(\mathbf{w}'_1, \dots, \mathbf{w}'_m)$, where $\mathbf{v}'_i = \sum_{j=1}^n q_{ji} \mathbf{v}_j$ and $\mathbf{w}'_i = \sum_{j=1}^m p_{ji} \mathbf{w}_j$, or

$$\boxed{\vec{v}' = Q^T \vec{v} \text{ and } \vec{w}' = P^T \vec{w}}.$$

There is a corresponding change of coordinates:

$$\mathbf{x} = \vec{v}'^T \vec{x}' = \vec{v}^T Q \vec{x}' \quad \text{so} \quad \boxed{\vec{x} = Q \vec{x}'}$$

Similarly,

$$\boxed{\vec{b} = P \vec{b}'} \quad \text{and} \quad \boxed{A = P A' Q^{-1}}.$$

E. Matrix normal forms in terms of linear transformations

Again look at $L : V \rightarrow W$. Consider for the moment a basis change only in W : $\vec{w}' = P^T \vec{w}$. Then the matrices A and A' representing L are related by $A = PA'$, that is, they are *row equivalent*. Thus we have the

Theorem: *Given a linear mapping $L : V \rightarrow W$ and a fixed basis \vec{v} for V there is a choice of basis \vec{w} for W such that the matrix representing L in these bases is in RREF. Moreover, that matrix is uniquely determined by L .*

- This is the second way in which normal forms arise in linear algebra: we have a (fixed) abstract object, such as a linear transformation, and we want to choose a basis (or bases) in which the matrix representing it has a standard form.
- Our previous formulation may be looked at as a special case of this one: take $V = \mathbb{F}^n$ and $W = \mathbb{F}^m$, and identify $L : V \rightarrow W$ with the matrix A such that $T\vec{x} = A\vec{x}$.

II. Normal forms for $L : V \rightarrow W$ or for an $m \times n$ matrix

Equivalence relation on matrices (P, Q invertible)	Basis choice in vector space	Normal form	Invariants
Row equivalence $A \sim PA$	Choice of basis in W	Reduced row echelon form	r, j_1, \dots, j_r • entries
Column equivalence $A \sim AQ$	Choice of basis in V	Reduced column echelon form	r, i_1, \dots, i_r • entries
Equivalence $A \sim PAQ$	Choice of basis in both V and W	$\begin{bmatrix} I_r & 0 \\ 0 & 0 \end{bmatrix}$	r

III. Normal forms for $L : V \rightarrow V$

- V an n dimensional vector space, basis $\vec{v}^T = [\mathbf{v}_1, \dots, \mathbf{v}_n]$
- $L : V \rightarrow V$ linear with matrix A in this basis: $L\mathbf{x} = \vec{v}^T A \vec{x}$.

Problem: Find a normal form for the matrix representing L in a suitable basis.

- The matrix representing L in a new basis $\vec{v}' = Q^T \vec{v}$ is then $A' = Q^{-1} A Q$. We say that A and A' are *similar*.

Equivalent problem: Find a normal form for $n \times n$ matrices under the equivalence relation of similarity.

- In this section we will have to pay attention to the field \mathbb{F} over which our spaces are defined. We will always take \mathbb{F} to be either \mathbb{R} or \mathbb{C} .

A. Eigenvalues and eigenvectors

- A *nonzero* vector $\mathbf{x} \in V$ is an *eigenvector* of L , with *eigenvalue* λ , if $L\mathbf{x} = \lambda\mathbf{x}$.
- Suppose we can find a basis $(\mathbf{v}_1, \dots, \mathbf{v}_n)$ of V such that each \mathbf{v}_i is an eigenvector of L : $L\mathbf{v}_i = \lambda_i\mathbf{v}_i$. Then the matrix A of L in this basis is *diagonal*:

$$A = \begin{bmatrix} \lambda_1 & 0 & \cdots & 0 \\ 0 & \lambda_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \lambda_n \end{bmatrix}.$$

We say that L has been *diagonalized*. Thus if we could always find such a basis our normal form would be a diagonal matrix. But in fact we cannot do so in general.

- In the alternate language: let A be an $n \times n$ matrix. We ask if A can be *diagonalized*, i.e., if there exists an matrix Q such that $A' = Q^{-1} A Q$ is diagonal. But *not all matrices can be diagonalized*.

Diagonal matrices are so pleasant to encounter that we add a new problem to our list:

Problem: Find special classes of linear transformations such that every member of the class can be diagonalized.

Remark on uniqueness: Recall that the eigenvalues of a matrix A are precisely the roots of the characteristic polynomial

$$p_A(\lambda) = \det(A - \lambda I).$$

Moreover, if A is diagonalizable then the number of times each eigenvalue occurs in the diagonal form is equal to its multiplicity as a root of the characteristic polynomial. Thus the diagonal form is *unique up to a rearrangement of the diagonal entries*.

- This is an example of what we referred to earlier as an “essentially unique” normal form.
- If $\mathbb{F} = \mathbb{C}$ then L always has at least one eigenvalue; *this may not be true if $\mathbb{F} = \mathbb{R}$.*

B. Normal form for general $L : V \rightarrow V$ over \mathbb{C}

Theorem: (a) Suppose that V is a vector space over \mathbb{C} , the field of complex numbers. Then every linear transformation $L : V \rightarrow V$ is represented in an appropriately chosen basis by an essentially unique matrix in *Jordan canonical form*.

(b) Equivalently, every complex matrix is similar to a matrix in Jordan canonical form.

Idea: If λ is an eigenvalue of L then we may define the *eigenspace*

$$V_\lambda = \{ \mathbf{x} \in V \mid (L - \lambda I)\mathbf{x} = 0 \}.$$

L is diagonalizable if and only if $V = \bigoplus_\lambda V_\lambda$, that is, if and only if every $\mathbf{x} \in V$ can be written as a sum $\mathbf{x} = \sum_\lambda \mathbf{x}_\lambda$, with $\mathbf{x}_\lambda \in V_\lambda$.

Remark: The notation $V = \bigoplus_\lambda V_\lambda$ means that V is the *direct sum* of the subspaces V_λ . In fact, this requires that every $\mathbf{x} \in V$ can be written *uniquely* as a sum $\mathbf{x} = \sum_\lambda \mathbf{x}_\lambda$, with $\mathbf{x}_\lambda \in V_\lambda$. When the V_λ are the eigenspaces, however, this uniqueness is automatically satisfied.

We may also define the *generalized eigenspace*

$$\widehat{V}_\lambda = \{ \mathbf{x} \in V \mid (L - \lambda I)^k \mathbf{x} = 0 \text{ for some } k \in \mathbb{N} \}.$$

Now always $V = \bigoplus_\lambda \widehat{V}_\lambda$; moreover, \widehat{V}_λ is *invariant* under L : $L(\widehat{V}_\lambda) \subset \widehat{V}_\lambda$. So it suffices to study the restriction of L to \widehat{V}_λ .

We know that for each $\mathbf{x} \in \widehat{V}_\lambda$ there is an integer k , which *a priori* may depend on \mathbf{x} , such that $(L - \lambda I)^k \mathbf{x} = 0$. Because V is finite dimensional, however, there must be some one k which has this property for *all* $\mathbf{x} \in \widehat{V}_\lambda$. That is, when restricted to \widehat{V}_λ , L satisfies $(L - \lambda I)^k = 0$ for some k . How could this happen?

Suppose there are vectors $\mathbf{v}_1^*, \dots, \mathbf{v}_j^* \in \widehat{V}_\lambda$, with $j \leq k$, so that

$$(L - \lambda I)\mathbf{v}_j^* = \mathbf{v}_{j-1}^*, \quad (L - \lambda I)\mathbf{v}_{j-1}^* = \mathbf{v}_{j-2}^*, \quad \dots \quad (L - \lambda I)\mathbf{v}_1^* = 0.$$

Then $(L - \lambda I)^j \mathbf{v}_i^* = 0 = (L - \lambda I)^k \mathbf{v}_i^*$ for any i . Let V^* be the subspace of V spanned by these vectors; in that subspace, and using $\mathbf{v}_1^*, \dots, \mathbf{v}_j^*$ as a basis, L has matrix

$$J^\lambda = \begin{bmatrix} \lambda & 1 & 0 & \cdots & 0 & 0 \\ 0 & \lambda & 1 & \cdots & 0 & 0 \\ 0 & 0 & \lambda & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & 1 & 0 \\ 0 & 0 & 0 & \cdots & \lambda & 1 \\ 0 & 0 & 0 & \cdots & 0 & \lambda \end{bmatrix}.$$

A matrix of this form is called a *Jordan block*.

Fact: We can always find a basis of \widehat{V} consisting of one or more sequences $\mathbf{v}_1^*, \dots, \mathbf{v}_j^*$ of vectors of the above type.

Let L have *distinct* eigenvalues $\lambda_1, \dots, \lambda_m$ then we may find a basis for \widehat{V}_{λ_i} in which the restriction of L to \widehat{V}_{λ_i} has the *block diagonal* form

$$K_i = \begin{bmatrix} J_1^{\lambda_1} & & \\ & \cdots & \\ & & J_{r_1}^{\lambda_1} \end{bmatrix}, \quad (*)$$

with each $J_r^{\lambda_1}$ a Jordan block. In the basis for V which is the union of all of these bases, L has the form

$$A = \begin{bmatrix} K_1 & & \\ & \cdots & \\ & & K_m \end{bmatrix}. \quad (**)$$

This is the Jordan form for L .

Uniqueness: L uniquely determines the *number* and *sizes* of the Jordan blocks associated with each eigenvalue. (These, and the eigenvalues themselves, are the invariants.) One may by convention group together all blocks associated with the same eigenvalue, as in (**), and arrange the blocks in each K_i (see (*)) in order of (say) decreasing size. However, the order of the K_i in (**) is not unique.

Remark: L satisfies its own characteristic equation, i.e., $p_L(L) = 0$. The monic polynomial $q_L(\lambda)$ of lowest degree for which $q_L(L) = 0$ is called the *minimal polynomial* of L ; q_L divides p_L . In fact, if k_i is the size of the matrix K_i —the dimension of \widehat{V}_{λ_i} —and ℓ_i is the size of the largest Jordan block J^{λ_i} appearing in the Jordan form of L , then

$$p_L(\lambda) = \prod_{i=1}^m (\lambda - \lambda_i)^{k_i} \quad \text{and} \quad q_L(\lambda) = \prod_{i=1}^m (\lambda - \lambda_i)^{\ell_i}.$$

Example: if the Jordan form is (all blank entries are 0)

$$\left[\begin{array}{ccc|cc} 3 & 1 & & & \\ & 3 & 1 & & \\ & & 3 & & \\ \hline & & & 3 & 1 \\ & & & & 3 \\ & & & & & 3 \\ & & & & & & 0 \\ & & & & & & & 0 \end{array} \right]$$

then $p(\lambda) = \lambda^2(\lambda - 3)^6$ and $q(\lambda) = \lambda(\lambda - 3)^3$.

C. Computation with matrices

The Jordan form is useful for computing powers and exponentials of matrices. For example, suppose we want to solve a linear system of differential equations for an unknown vector function $\vec{x}(t)$:

$$\frac{d\vec{x}}{dt} = A \vec{x}, \quad \vec{x}(0) = \vec{x}_0,$$

where A (an $n \times n$ matrix) and \vec{x}_0 are given. Formally the solution is $\vec{x}(t) = e^{tA}\vec{x}_0$, and the matrix exponential may be defined by a power series, but we would like to have an effective way to compute it.

Suppose $A = Q^{-1}KQ$ with K in Jordan form; then

$$e^{tA} = \sum_{k=0}^{\infty} \frac{t^k}{k!} A^k = \sum_{k=0}^{\infty} \frac{t^k}{k!} (Q^{-1}KQ)^k = \sum_{k=0}^{\infty} \frac{t^k}{k!} Q^{-1} K^k Q = Q^{-1} e^{tK} Q,$$

so we need only compute e^{tK} .

But

$$K = \begin{bmatrix} J_1 & & \\ & \cdots & \\ & & J_r \end{bmatrix} \implies e^{tK} = \begin{bmatrix} e^{tJ_1} & & \\ & \cdots & \\ & & e^{tJ_r} \end{bmatrix}$$

so we need only to compute the exponential of a single Jordan block.

This turns out to be an easy exercise. Suppose for example that J is an $n \times n$ Jordan block with eigenvalue λ , then

$$e^{tJ} = e^{t\lambda} \begin{bmatrix} 1 & t & \frac{t^2}{2!} & \cdots & \frac{t^{n-1}}{(n-1)!} \\ 0 & 1 & t & \cdots & \frac{t^{n-2}}{(n-2)!} \\ 0 & 0 & 1 & \cdots & \frac{t^{n-3}}{(n-3)!} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & 1 \end{bmatrix}$$

D. Normal form for general $L : V \rightarrow V$ over \mathbb{R}

Here the normal form is *real canonical form*; this is less important than the Jordan canonical form, and we will be brief. Some eigenvalues of L may be real; others come in complex conjugate pairs $\alpha \pm i\beta$. The normal form again has blocks along the diagonal:

- Corresponding to the real eigenvalues there are Jordan blocks;
- Corresponding to the eigenvalues $\alpha \pm i\beta$ there are blocks

$$\begin{bmatrix} \Lambda & I & & & & & \\ & \Lambda & I & & & & \\ & & \Lambda & \cdots & & & \\ & & & \cdots & I & & \\ & & & & \Lambda & I & \\ & & & & & \Lambda & \\ & & & & & & \Lambda \end{bmatrix}.$$

where

$$I = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \Lambda = \begin{bmatrix} \alpha & \beta \\ -\beta & \alpha \end{bmatrix}.$$

E. When is $L : V \rightarrow V$, or a matrix A , necessarily diagonalizable?

Let us first consider matrices. We recall some terminology:

- The *adjoint* A^* of a complex matrix A is the transpose of its complex conjugate: $A^* = \overline{A}^T$. A is *symmetric* if $A^T = A$ and *self-adjoint* or *Hermitian symmetric* if $A^* = A$.
- A real matrix P is *orthogonal* if $P^T = P^{-1}$, and a complex matrix P is *unitary* if $P^* = P^{-1}$. We say that a matrix A is *diagonalizable by an orthogonal (or unitary) matrix* if $P^{-1}AP$ is diagonal for some orthogonal (unitary) P .

Theorem: (a) *If A is a real symmetric $n \times n$ matrix then all the eigenvalues of A are real and A is diagonalizable by an orthogonal matrix.*

(b) *If A is a self-adjoint $n \times n$ matrix then all the eigenvalues of A are real and A is diagonalizable by a unitary matrix.*

To state the corresponding result for $L : V \rightarrow V$ we assume that V is an *inner product space*.

- Recall: for $\mathbf{x}, \mathbf{y} \in V$ the inner product $\langle \mathbf{x}, \mathbf{y} \rangle$ lies in \mathbb{F} ; the inner product is *linear in its first argument*, *positive definite* ($\langle \mathbf{x}, \mathbf{x} \rangle \geq 0$, with equality only if $\mathbf{x} = 0$), and *symmetric* ($\langle \mathbf{x}, \mathbf{y} \rangle = \langle \mathbf{y}, \mathbf{x} \rangle$) if $\mathbb{F} = \mathbb{R}$, *Hermitian symmetric* ($\langle \mathbf{x}, \mathbf{y} \rangle = \overline{\langle \mathbf{y}, \mathbf{x} \rangle}$) if $\mathbb{F} = \mathbb{C}$.
- Note that when $\mathbb{F} = \mathbb{C}$ the inner product is *antilinear* in its second argument: $\langle \mathbf{x}, a\mathbf{y} + b\mathbf{z} \rangle = \bar{a}\langle \mathbf{x}, \mathbf{y} \rangle + \bar{b}\langle \mathbf{x}, \mathbf{z} \rangle$.
- A basis for $\mathbf{v}_1, \dots, \mathbf{v}_n$ for V is *orthonormal* if $\langle \mathbf{v}_i, \mathbf{v}_j \rangle = \delta_{ij}$.
- If $\vec{\mathbf{v}}$ and $\vec{\mathbf{v}}' = P^T \vec{\mathbf{v}}$ are two orthonormal bases then P is (i) orthogonal if $\mathbb{F} = \mathbb{R}$ and (ii) unitary if $\mathbb{F} = \mathbb{C}$.
- If $\langle L\mathbf{x}, \mathbf{y} \rangle = \langle \mathbf{x}, L\mathbf{y} \rangle$ for all $\mathbf{x}, \mathbf{y} \in V$ then L is *symmetric* when $\mathbb{F} = \mathbb{R}$ and *self-adjoint* or *Hermitian symmetric* when $\mathbb{F} = \mathbb{C}$.
- If $\vec{\mathbf{v}}$ is an orthonormal basis then the matrix A of L in this basis is (i) symmetric if $\mathbb{F} = \mathbb{R}$ and (ii) self-adjoint if $\mathbb{F} = \mathbb{C}$.

Thus the previous theorem is equivalent to

Theorem: *Suppose that either $\mathbb{F} = \mathbb{R}$ and L is symmetric, or that $\mathbb{F} = \mathbb{C}$ and L is self-adjoint. Then all the eigenvalues of L are real and there exists an orthonormal basis for V composed of eigenvectors of L .*

We will sketch the proof of this theorem for the case $\mathbb{F} = \mathbb{C}$.

- First we show that all the eigenvalues of L are real. If $L\mathbf{x} = \lambda\mathbf{x}$ then

$$\begin{aligned}\langle L\mathbf{x}, \mathbf{x} \rangle &= \langle \lambda\mathbf{x}, \mathbf{x} \rangle = \lambda\langle \mathbf{x}, \mathbf{x} \rangle \\ &= \langle \mathbf{x}, L\mathbf{x} \rangle = \langle \mathbf{x}, \lambda\mathbf{x} \rangle = \bar{\lambda}\langle \mathbf{x}, \mathbf{x} \rangle\end{aligned}$$

Thus $\lambda\langle \mathbf{x}, \mathbf{x} \rangle = \bar{\lambda}\langle \mathbf{x}, \mathbf{x} \rangle$ and since $\langle \mathbf{x}, \mathbf{x} \rangle \neq 0$, $\lambda = \bar{\lambda}$.

- Now we use induction on the dimension of V to find an orthonormal basis of eigenvectors. Such a basis obviously exists if $\dim V = 1$. If $\dim V = d > 1$ then we know that L has at least one eigenvalue and corresponding eigenvector, say $L\mathbf{x} = \lambda\mathbf{x}$. Let W be the subspace of V consisting of all vectors orthogonal to \mathbf{x} :

$$W = \{ \mathbf{y} \in V \mid \langle \mathbf{x}, \mathbf{y} \rangle = 0 \}.$$

If $\mathbf{y} \in W$ then $L\mathbf{y} \in W$, since

$$\langle \mathbf{x}, L\mathbf{y} \rangle = \langle L\mathbf{x}, \mathbf{y} \rangle = \langle \lambda\mathbf{x}, \mathbf{y} \rangle = \lambda \langle \mathbf{x}, \mathbf{y} \rangle = 0.$$

- Thus the restriction L_W of L to W is a mapping $L_W : W \rightarrow W$; L_W is easily seen to be self-adjoint and thus there is an orthonormal basis for W consisting of eigenvectors of L_W . Adding \mathbf{x} to this basis gives us the require basis of V . ■

There is an important consequence. A (complex) matrix A is *normal* if A and its adjoint A^* commute: $AA^* = A^*A$.

Corollary: *If A is a normal $n \times n$ matrix then A is diagonalizable by a unitary matrix.*

Proof sketch: Let $A_R = (A + A^*)/2$ and $A_I = (A - A^*)/2i$; it is easy to see that A_R and A_I are self-adjoint and that $A = A_R + iA_I$. Because A and A^* commute, so do A_R and A_I . We know that we may diagonalize each of these separately with a unitary matrix, but *because they commute we can find a single unitary matrix that diagonalizes both*, and that is all we need.

To understand the key step here we may first diagonalize A_R . Let V_λ be an eigenspace of A_R ; from $A_RA_I = A_IA_R$ we find that A_I maps V_λ into itself, and therefore can be diagonalized within V_λ . Putting together these diagonalizations in each eigenspace of A_R gives us the joint diagonalization of A_R and A_I .

IV. Bilinear and quadratic forms

Let V be a finite dimensional vector space; for the moment take $\mathbb{F} = \mathbb{R}$.

- A *bilinear form* B on V is a mapping $B : V \times V \rightarrow \mathbb{F}$ which is linear in each argument:

$$B(ax + by, z) = aB(x, z) + bB(y, z)$$

$$B(x, ay + bz) = aB(x, y) + bB(x, z)$$

B is *symmetric* if $B(x, y) = B(y, x)$ for all $x, y \in V$ and *skew symmetric* or *alternating* if $B(x, y) = -B(y, x)$ for all x, y .

- A *quadratic form* Q on V is a mapping $Q : V \rightarrow \mathbb{F}$ obtained from a symmetric bilinear form B via $Q(x) = B(x, x)$.
- **Remark:** one may give an intrinsic definition of a quadratic form and then show that, *because \mathbb{F} does not have characteristic 2*, quadratic forms and symmetric bilinear forms are equivalent in the sense above. Over a field of characteristic 2 these are not equivalent concepts.

- If B is a bilinear form on V and $\mathbf{v}_1, \dots, \mathbf{v}_n$ is a basis for V then one may associate with B the matrix A defined by $a_{ij} = B(\mathbf{v}_i, \mathbf{v}_j)$.
- If \mathbf{x} and \mathbf{y} are vectors in V and \vec{x} and \vec{y} their coordinate vectors ($\mathbf{x} = \vec{v}^T \vec{x}$, $\mathbf{y} = \vec{v}^T \vec{y}$) then $B(\mathbf{x}, \mathbf{y}) = \vec{x}^T A \vec{y}$. B is symmetric if $A = A^T$ and skew symmetric if $A = -A^T$.
- Under a change of basis $\vec{v}' = P^T \vec{v}$ we have $\vec{x} = P \vec{x}'$, $\vec{y} = P \vec{y}'$, and so $B(\mathbf{x}, \mathbf{y}) = \vec{x}^T A \vec{y} = \vec{x}'^T P^T A P \vec{y}'$. Thus

$$\boxed{A' = P^T A P} \quad (*)$$

When A and A' are related as in $(*)$, with P an invertible matrix, we say that they are *congruent*.

Problem: Find normal forms for real symmetric matrices, and for real skew symmetric matrices, under the equivalence relation of congruence.

Equivalent problem: Find normal forms for symmetric and skew symmetric bilinear forms, on a real vector space, under change of basis.

A. Symmetric bilinear forms

- Let A be a real symmetric matrix. We know that A can be diagonalized by an orthogonal matrix P :

$$P^{-1}AP = \begin{bmatrix} \lambda_1 & & \\ & \cdots & \\ & & \lambda_n \end{bmatrix}.$$

Since $P^{-1} = P^T$, this implies that every symmetric matrix is congruent, via an orthogonal matrix, to a diagonal matrix. The matrix entries on the diagonal are of course the eigenvalues of A .

- Equivalently: if B is a bilinear form on a real vector space V then there is an orthonormal basis \vec{v} for V such that for any $\mathbf{x}, \mathbf{y} \in V$

$$B(\mathbf{x}, \mathbf{y}) = \sum_{i=1}^n \lambda_i x_i y_i,$$

where as usual $\mathbf{x} = \vec{v}^T \vec{x}$ and $\mathbf{y} = \vec{v}^T \vec{y}$.

B. An application to mechanics

Consider a system of n particles

- coordinates $x_1(t), \dots, x_n(t)$, velocities $\dot{x}_1(t), \dots, \dot{x}_n(t)$.
- kinetic energy $T = \dot{\vec{x}}^T M \dot{\vec{x}}$, potential energy $V = \vec{x}^T Q \vec{x}$, with M and Q symmetric matrices and M positive definite (all eigenvalues strictly positive).

- Equations of motion:
$$M \ddot{\vec{x}} = -Q \vec{x}$$

n *coupled* equations in n unknowns.

Idea: Introduce new coordinates \vec{y} with $\vec{x} = P\vec{y}$ (and hence $\dot{\vec{x}} = P\dot{\vec{y}}$) so that in the new coordinates the kinetic and potential energies are given by

$$T = \dot{\vec{y}}^T P^T M P \dot{\vec{y}} = \dot{\vec{y}}^T M' \dot{\vec{y}}, \quad V = \vec{y}^T P^T Q P \vec{y} = \vec{y}^T Q' \vec{y},$$

with **diagonal** matrices $M' = P^T M P$ and $Q' = P^T Q P$.

How?

- First find P_1 with $P_1^T M P_1 = I$ (possible because M is positive definite). Q is now replaced by $Q_1 = P_1^T Q P_1$.

- Now find an **orthogonal** P_2 with $Q' = P_2^T Q_1 P_1 = \begin{bmatrix} \lambda_1 & & \\ & \cdots & \\ & & \lambda_n \end{bmatrix}$ diagonal, and take $P = P_2 P_1$, $\vec{x} = P \vec{y}$. Then

$$T = \dot{\vec{y}}^T P_2^T I P_2 \dot{\vec{y}} = \dot{\vec{y}}^T \dot{\vec{y}} = \sum_i \dot{y}_i^2, \quad V = \vec{y}^T Q' \vec{y} = \sum_i \lambda_i y_i^2$$

- Now the equations of motion decouple: for $i = 1, \dots, n$, if $\lambda_i < 0$,

$$\ddot{y}_i = -\lambda_i y_i \quad \implies \quad y_i(t) = A_i \cos \omega_i t + B_i \sin \omega_i t$$

for $\omega_i = \sqrt{-\lambda_i}$. Thus

$$\vec{x}_j(t) = \sum_{i=1}^n p_{ij} (A_i \cos \omega_i t + B_i \sin \omega_i t)$$

The motion of \vec{x} is a superposition of “normal modes” of vibration.

C. Skew symmetric bilinear forms

Theorem: *Every real $n \times n$ skew symmetric matrix A is congruent to a unique matrix with a block decomposition of the form*

$$\begin{bmatrix} 0 & I & 0 \\ -I & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Here I is the identity matrix of some dimension r with $2r \leq n$.